

# Recent Developments in Event Generators

Stefan Höche

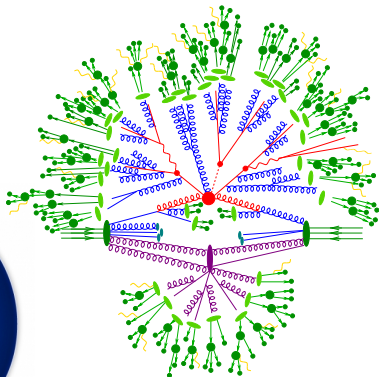
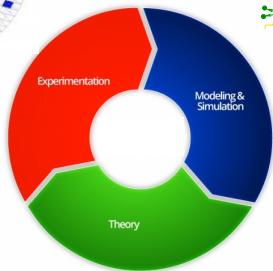
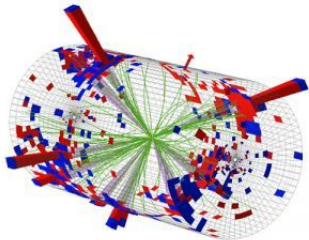
Fermi National Accelerator Laboratory

QCD@LHC

IJClab Saclay, 29/11/2022



# The toolkit



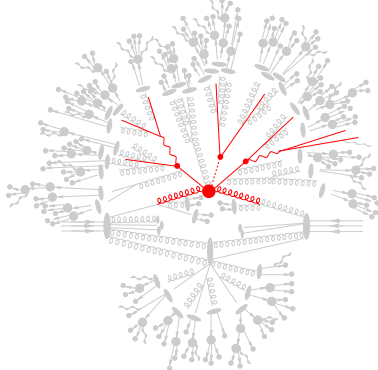
$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + h.c.$$

# LHC event generators

[Buckley et al.] arXiv:1101.2599

[Campbell et al.] arXiv:2203.11110

- ▶ Short distance interactions
  - ▶ **Signal process**
  - ▶ Radiative corrections
- ▶ Long-distance interactions
  - ▶ Hadronization
  - ▶ Particle decays



## Divide and Conquer

- ▶ Quantity of interest: Total interaction rate
- ▶ Convolution of short & long distance physics

$$\sigma_{p_1 p_2 \rightarrow X} = \sum_{i,j \in \{q,g\}} \int dx_1 dx_2 \underbrace{f_{p_1,i}(x_1, \mu_F^2) f_{p_2,j}(x_2, \mu_F^2)}_{\text{long distance}} \underbrace{\hat{\sigma}_{ij \rightarrow X}(x_1 x_2, \mu_F^2)}_{\text{short distance}}$$

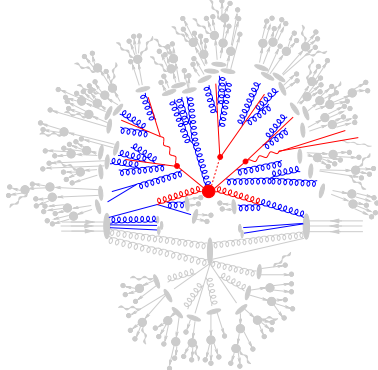


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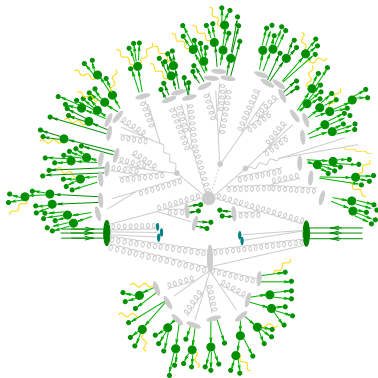
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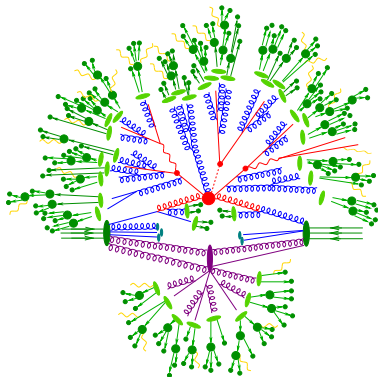
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## Connection to QCD theory

- ▶  $\hat{\sigma}_{ij \rightarrow n}(\mu_F^2) \rightarrow$  Collinearly factorized fixed-order result at N<sup>x</sup>LO

Implemented in fully differential form to be maximally useful

Tree level:  $d\Phi_n B_n$

- ▶ Automated ME generators + phase-space integrators

1-Loop level:  $d\Phi_n \left( B_n + V_n + \sum C + \sum I_n \right) + d\Phi_{n+1} \left( R_n - \sum S_n \right)$

- ▶ Automated loop ME generators + integral libraries + IR subtraction

2-Loop level: It depends ...

- ▶ Individual solutions based on SCET,  $q_T$  subtraction, P2B

- ▶  $f_i(x, \mu_F^2) \rightarrow$  Collinearly factorized PDF at N<sup>y</sup>LO

Evaluated at  $O(1\text{GeV}^2)$  and expanded into a series above  $1\text{GeV}^2$

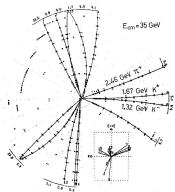
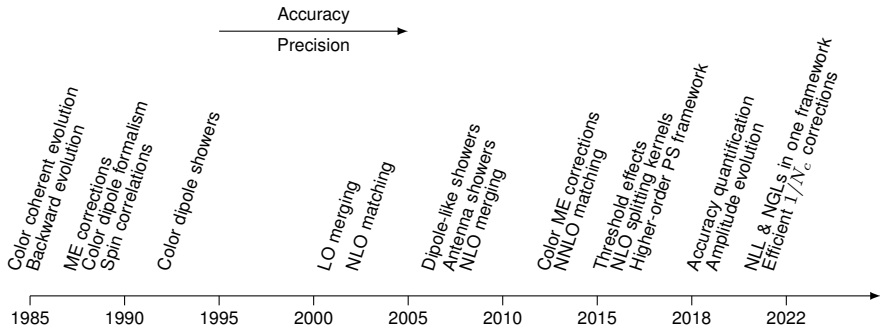
$$\text{DGLAP: } \frac{dx x f_a(x, t)}{d \ln t} = \sum_{b=q,g} \int_0^1 d\tau \int_0^1 dz \frac{\alpha_s}{2\pi} [z P_{ab}(z)]_+ \tau f_b(\tau, t) \delta(x - \tau z)$$

- ▶ Parton showers, dipole showers, antenna showers, ...

$$\text{Matching: } d\Phi_n \frac{S_n}{B_n} \leftrightarrow \frac{dt}{t} dz \frac{\alpha_s}{2\pi} P_{ab}(z)$$

- ▶ MC@NLO, POWHEG, Geneva, MINNLO<sub>PS</sub>, ...

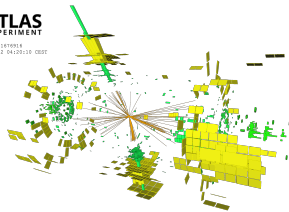
# Co-design of simulations over the years



$\sqrt{s} \times 500$   
→  
 $e^+e^-$  vs.  $pp$

**ATLAS**  
EXPERIMENT

Event: 331670918  
2013-09-22 04:10:10 CERN



# Directions of development

## Much effort focused on parton-shower component recently

- ▶ Phenomenologically interesting: Drives jet production,  $b$ -tagging, ...
- ▶ Experimentally relevant: Often source of largest uncertainty
- ▶ Next to hadronization, probably the most important component of MCs

## Fixed-order aspects

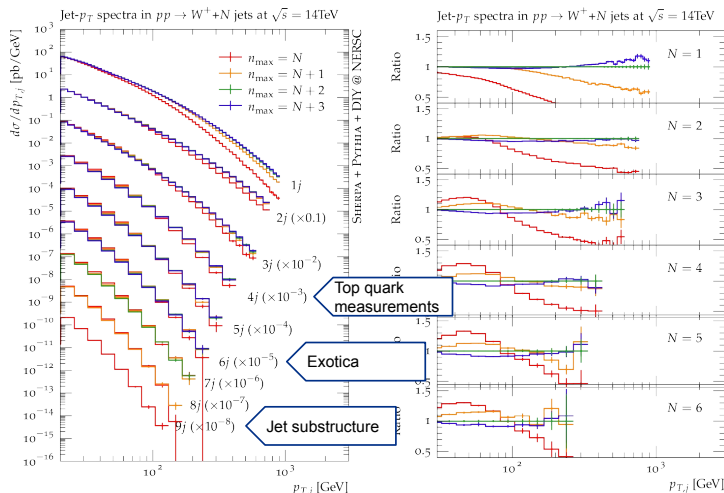
- ▶ Matching to NLO & merging
  - ▶ Negative weight fraction
  - ▶ Computing efficiency
- ▶ Matching to NNLO calculations
  - ▶ Semi-inclusive (Geneva, MINNLO<sub>PS</sub>)
  - ▶ Fully differential (Vincia)
- ▶ Matching to N<sup>3</sup>LO calculations
  - ▶ Fully differential (TOMTE)

## All-order aspects

- ▶ NLL precision
- ▶ NLO splitting functions
- ▶ Kinematic edge effects
- ▶ Spin correlations in collinear & soft limit
- ▶ Sub-leading color effects
- ▶ Threshold effects
- ▶ Amplitude evolution

# Why matching & merging?

[Prestel,Schulz,SH] arXiv:1905.05120



- ▶ Predictions for measured  $N$ -jet rates stabilize for  $\approx N+2$  LO ME-level jets
- ▶ Poor man's version of NNLO (loops emulated by legs + unitarity constraint)

# Computing efficiency: The cost of multi-jet merging

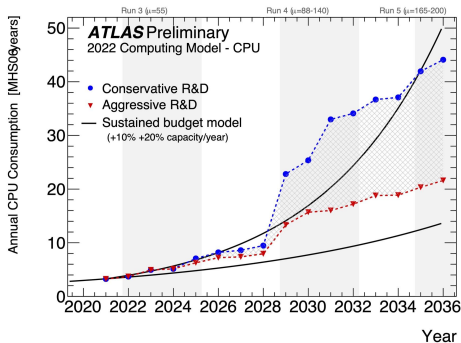
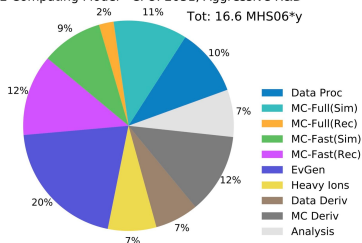
[HSF Generator WG] arXiv:2004.13687, arXiv:2109.14938

- ▶ Event generation will consume significant fraction of resources at LHC soon
- ▶ Need to scrutinize both generator usage and underlying algorithms
- ▶ Dedicated effort in HEP Software Foundation (HSF)

**ATLAS Preliminary**

2022 Computing Model - CPU: 2031, Aggressive R&D

Tot: 16.6 MHS06\*y



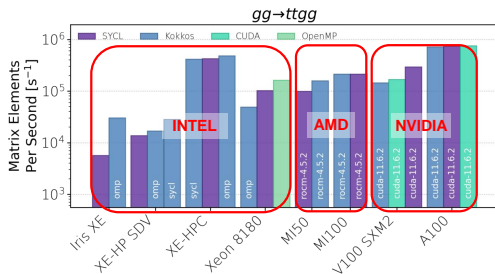
[ATLAS] CERN-LHCC-2022-005 / LHCC-G-182



# Computing efficiency: MadGraph Developments

[A. Valassi et al., ACAT'22 & QCD@LHC 2022]

- ▶ New code-generator in MadGraph 5 to generate CUDA, SYCL, Kokkos output for ME computation
- ▶ Vectorized code for computations on CPUs
- ▶ Included in improved MadEvent framework
- ▶ Performances of SYCL and Kokkos comparable to direct CUDA
- ▶ New computing strategy delivers both portability and performance

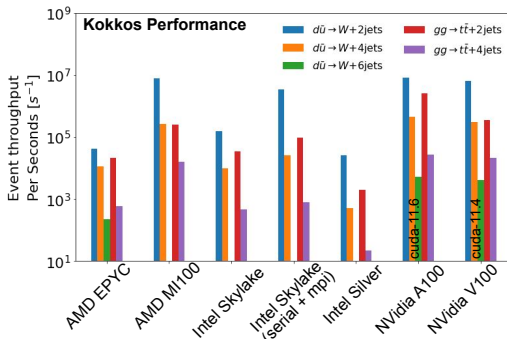


CUDA grid size		ACAT2022		madevent		standalone	
				8192		16384	
gg → t̄tggg	MEs precision	t <sub>TOT</sub> = t <sub>Mad</sub> + t <sub>MEs</sub> [sec]		N <sub>events</sub> /t <sub>TOT</sub> [events/sec]		N <sub>events</sub> /t <sub>MEs</sub> [MEs/sec]	
Fortran	double	1228.2	5.0 + 1223.2	7.34E1 (≈1.0)	7.37E1 (≈1.0)	—	—
CUDA	double	19.6	7.4 + 12.1	4.61E3 (x63)	7.44E3 (x100)	9.10E3	9.51E3 (x129)
CUDA	float	11.7	6.2 + 5.4	7.73E3 (x105)	1.66E4 (x224)	1.68E4	2.41E4 (x326)
CUDA	mixed	16.5	7.0 + 9.6	5.45E3 (x74)	9.43E3 (x128)	1.10E4	1.19E4 (x161)

# Computing efficiency: Sherpa Developments

[R. Wang et al., ACAT'22]

- ▶ Study of a variety of algorithms & assessment of practicality for LHC background simulations
- ▶ First use of new color basis [Melia] arXiv:1509.03297 in a generator
- ▶ Cuda for benchmarks, portability through Kokkos
- ▶ Factor  $\sim 10$  speedup at low multiplicity, factor  $\sim 4$  at high multiplicity (fully loaded E5620 CPU (MPI) and V100 GPU)
- ▶ Currently being combined with integrator and event generation framework

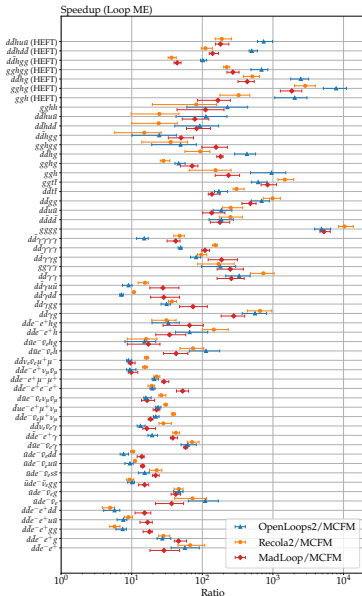


# Computing efficiency: Usage of analytics

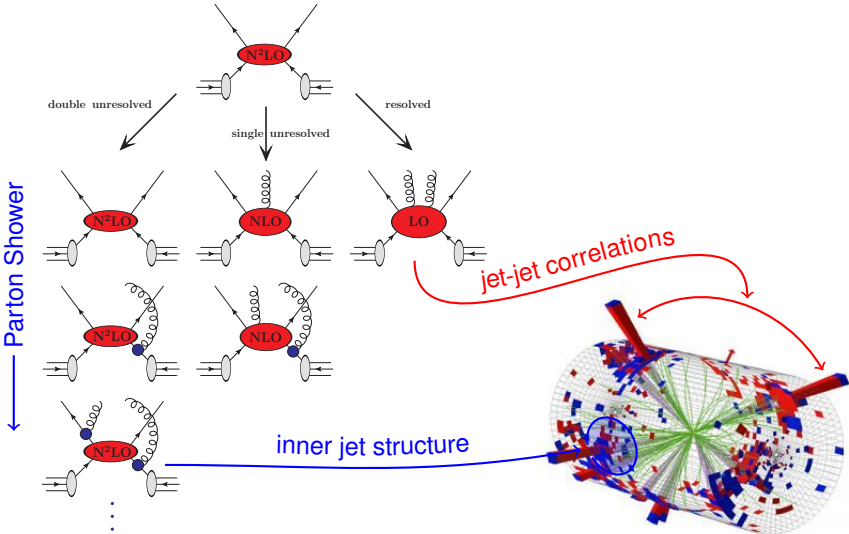
[Campbell,Preuss,SH] arXiv:2107.04472, [↗ M. Knobbe's talk]

- ▶ At HL-LHC, accuracy and precision requirements for a small number of processes drive computing demands:
  - ▶  $W^\pm/Z/\gamma$ +jets
  - ▶  $t\bar{t}$ +jets
  - ▶ ...
- ▶ Up to 2 jets, NLO matrix elements for  $W/Z/\gamma/h$  are known *analytically*
- ▶ Significant speedup out of the box (analytic vs numeric 1-loop ME only)

Merged Process	Sherpa+	Sherpa+
$n \leq 2$ @ NLO	OpenLoops2/MCFM	MadLoop5/MCFM
$n \leq 5$ @ LO		
$pp \rightarrow Z + n_j$	$1.83^{+0.20}_{-0.12}$	$3.01^{+0.26}_{-0.18}$
$pp \rightarrow W^+ + n_j$	$1.34^{+0.06}_{-0.07}$	$1.36^{+0.03}_{-0.03}$
$pp \rightarrow W^- + n_j$	$1.38^{+0.06}_{-0.04}$	$1.38^{+0.07}_{-0.11}$



# Fixed-order matching: Basic idea



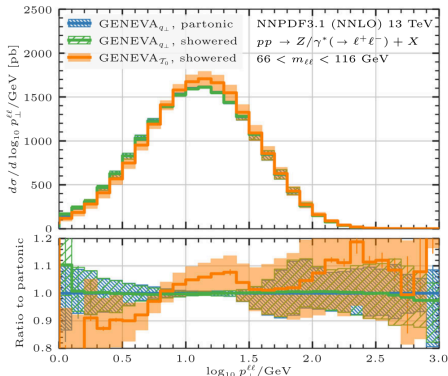
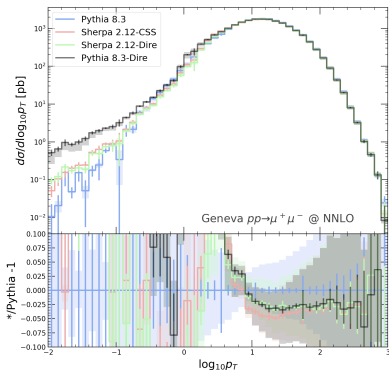
# Fixed-order matching: Geneva

[D. Napoletano's talk at HP<sup>2</sup>]

- Use known resummation in jettiness /  $q_T$  & match to NNLO

$$\frac{d\sigma}{d\Phi dr} = \frac{d\sigma^{\text{NNLL}'}}{d\Phi dr} - \frac{d\sigma^{\text{res.exp.}}}{d\Phi dr} + \frac{d\sigma^{\text{FO}}}{d\Phi dr}$$

- Match to shower by vetoing events with  $r_N(\Phi_{N+M}) > r_N$

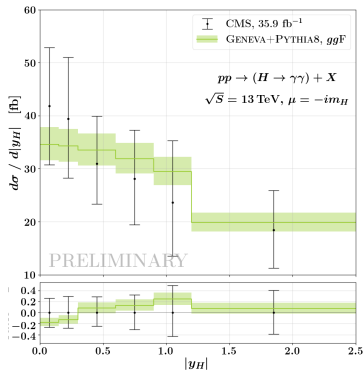
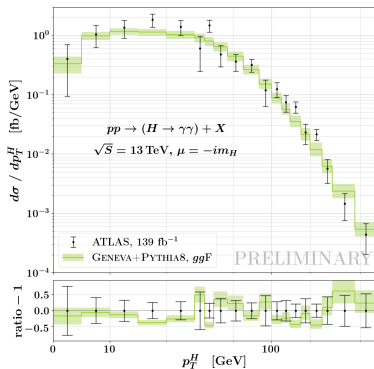


- Parton shower scheme uncertainty
- Choice of resolution variable

# Fixed-order matching: Geneva

[G. Marinelli's talk at HP<sup>2</sup>]

## ► Comparison against experimental data



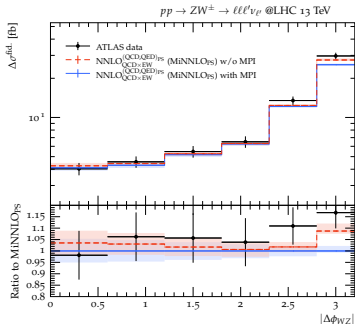
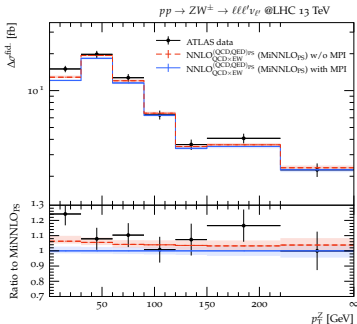
## ► $p_{T,H}$ and ATLAS data

## ► $y_H$ and CMS data

# Fixed-order matching: MINNLO<sub>PS</sub>

[Lindert,Lombardi,Wiesemann,Zanderighi,Zanoli] arXiv:2208.12660

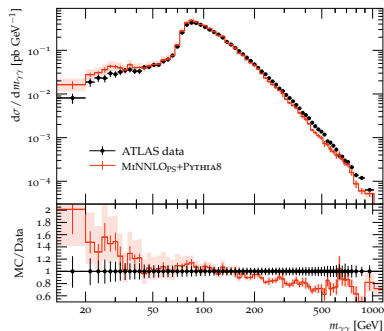
- ▶ WZ production at NNLO QCD  $\times$  NLO EW
- ▶ Various schemes to combine QCD & EW corrections  
→ associated uncertainty estimates



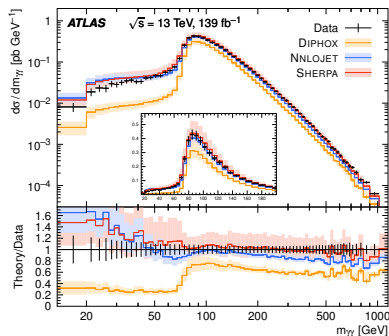
# Fixed-order matching: MINNLO<sub>PS</sub>

[Gavardi,Oleari,Re] arXiv:2204.12602

- ▶ Di-photon production at the LHC
- ▶ QED singular contributions in real-emission corrections treated as fixed order → split off by damping function



- ▶ Comparison between ATLAS data and MINNLO<sub>PS</sub>



- ▶ Previous experimental analysis [ATLAS] arXiv:2107.09330

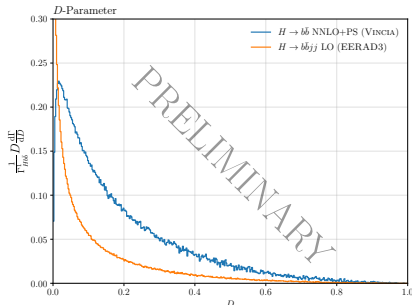
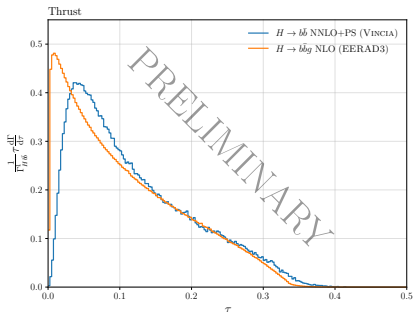


# Fixed-order matching: Vincia

[C. Preuss' talk at HP<sup>2</sup>]

[Campbell,Li,Preuss,Skands,SH] arXiv:2108.07133

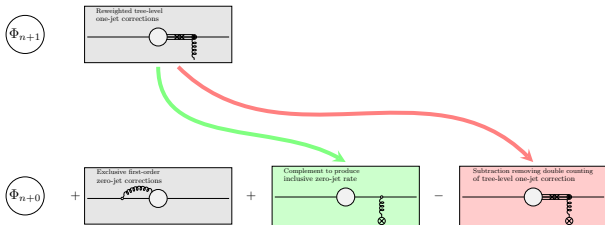
- ▶ Fully differential matching technique akin to POWHEG
- ▶ Technical implementation based on sector antenna framework
- ▶ Configurations absent in antenna-shower approximation simulated using direct  $2 \rightarrow 4$  branchings



# Fixed-order matching: $N^3\text{LO}$

[Lönnblad,Prestel] arXiv:1211.4827, [Plätzer] arXiv:1211.5467

## U(N)LOPS

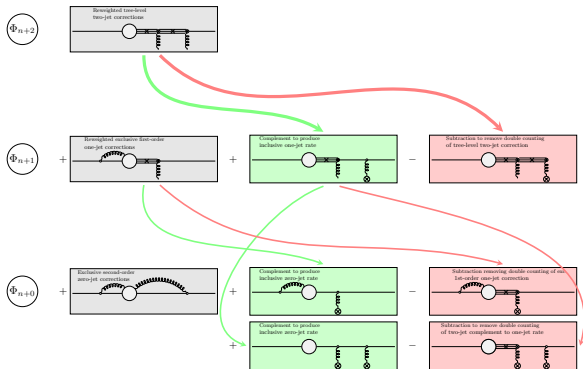


- ▶ Compute vetoed cross section & complete with real-emission
- ▶ Add Sudakov vetoed real-emission cross section & projection
- ▶ Can be implemented based on only two inputs (gray boxes)

# Fixed-order matching: $N^3\text{LO}$

[Lönnblad,Prestel] arXiv:1211.4827, [Li,Prestel,SH] arXiv:1405.3607

## $UN^2\text{LOPS}$

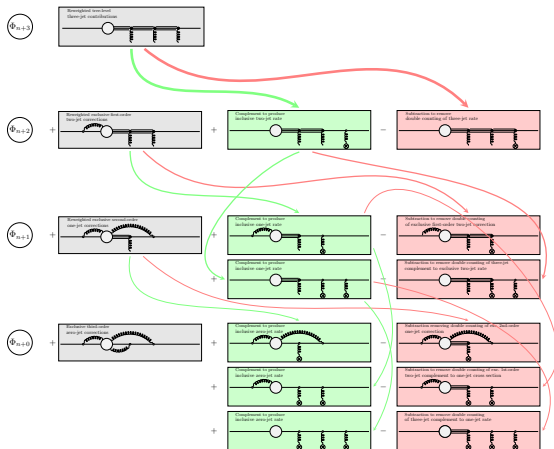


► Same idea as in ULOPS, but now also adding 2-loop contribution

# Fixed-order matching: $N^3LO$

[Prestel] arXiv:2106.03206, [Bertone,Prestel] arXiv:2202.01082

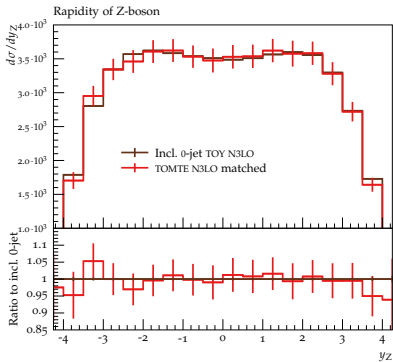
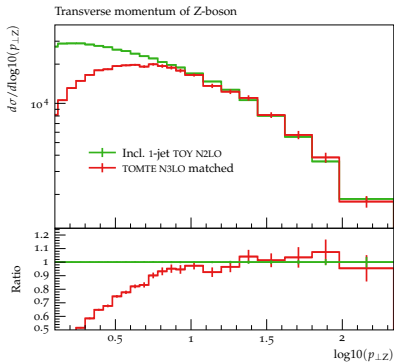
## TOMTE



- ▶ Same idea as in  $UN^2LOPS$ , but now also adding 3-loop contribution
- ▶ Must pay careful attention to projections (relevant for all  $UN^XLOPS$ )

# Fixed-order matching: $N^3\text{LO}$

[Bertone,Prestel] arXiv:2202.01082



- ▶ Drell-Yan lepton pair production at LHC
- ▶ Stand-in fixed-order calculation for closure tests

# All-order aspects: Parton showers at NLL precision

- ▶ How to quantify logarithmic precision of parton showers?  
[Dasgupta,Dreyer,Hamilton,Monni,Salam] arXiv:1805.09327
- ▶ Angular ordered parton showers provably NLL accurate for global observables, but wrong recoil may invalidate this  
[Bewick,Ferrario Ravasio,Richardson,Seymour] arXiv:1904.11866
- ▶ Two problems in commonly used dipole showers [[↗ talk by S. Ferrario-Ravasio](#)]
  - ▶ Correlations across multiple emissions due to recoil strategy
  - ▶ Color charge of initial quarks not reflected in soft, wide angle region
- ▶ Kinematics problem can be solved by
  - ▶ Partitioning of antenna radiation pattern, combined with local or semi-global recoil scheme [Dasgupta,Dreyer,Hamilton,Monni,Salam,Soyez] arXiv:2002.11114  
[vanBeekveld,Ferrario Ravasio,Hamilton,Salam,Soto-Ontoso,Soyez] arXiv:2205.02237, arXiv:2207.09467
  - ▶ Additive matching of soft to collinear radiator, combined with global recoil scheme [Forshaw,Holguin,Plätzer] arXiv:2003.06400
  - ▶ Multiplicative matching of soft to collinear radiator, combined with semi-global recoil scheme [Nagy,Soper] arXiv:2011.04773
  - ▶ Multiplicative matching of soft to collinear radiator, combined with global recoil scheme [Herren,Krauss,Reichelt,Schönherr,SH] arXiv:2208.06057



# Higher-order corrections: Collinear evolution at NLO

- Higher-order DGLAP evolution kernels obtained from factorization

$$D_{ji}^{(0)}(z, \mu) = \delta_{ij} \delta(1-z) \quad \leftrightarrow \quad \text{Diagram 1} / \text{Diagram 2}$$

Diagram 1: A grey circle with two incoming lines from the left and one outgoing line to the right labeled 'j' and 'z'.  
Diagram 2: A grey circle with two incoming lines from the left and one outgoing line to the right labeled 'i' and '1'.

$$D_{ji}^{(1)}(z, \mu) = -\frac{1}{\epsilon} P_{ji}^{(0)}(z) \quad \leftrightarrow \quad \text{Diagram 3} / \text{Diagram 4}$$

Diagram 3: A grey circle with two incoming lines from the left and one outgoing line to the right labeled 'j' and 'z'. A wavy line (gluon) is emitted from the top of the circle.  
Diagram 4: A grey circle with two incoming lines from the left and one outgoing line to the right labeled 'i' and '1'.

$$D_{ji}^{(2)}(z, \mu) = -\frac{1}{2\epsilon} P_{ji}^{(1)}(z) + \frac{\beta_0}{4\epsilon^2} P_{ji}^{(0)}(z) + \frac{1}{2\epsilon^2} \int_z^1 \frac{dx}{x} P_{jk}^{(0)}(x) P_{ki}^{(0)}(z/x)$$

$$\leftrightarrow \left( \text{Diagram 5} + \text{Diagram 6} \right) / \text{Diagram 4}$$

Diagram 5: Similar to Diagram 3, but the wavy line is emitted from the bottom of the circle.  
Diagram 6: Similar to Diagram 3, but the wavy line is emitted from the top of the circle and then splits into two wavy lines.

- $P_{ji}^{(n)}$  not probabilities, but sum rules hold ( $\leftrightarrow$  unitarity constraint)  
In particular: Momentum sum rule identical between LO & NLO
- Can perform the NLO computation of  $P_{ji}^{(1)}$  fully differentially using modified dipole subtraction [Catani, Seymour] hep-ph/9605323



# Higher-order corrections: Collinear evolution at NLO

[Prestel,SH] arXiv:1705.00742

- ▶ Example: Flavor-changing NLO splitting functions

$$P_{qq'}^{(1)}(z) = C_{qq'}(z) + I_{qq'}(z) + \int d\Phi_{+1} \left[ R_{qq'}(z, \Phi_{+1}) - S_{qq'}(z, \Phi_{+1}) \right]$$

- ▶ Real correction  $R_{qq'}$  and subtraction terms  $S_{qq'}$   
Difference finite in 4 dimensions  $\rightarrow$  amenable to MC simulation
- ▶ Integrated subtraction term and factorization counterterm given by

$$I_{qq'}(z) = \int d\Phi_{+1} S_{qq'}(z, \Phi_{+1})$$

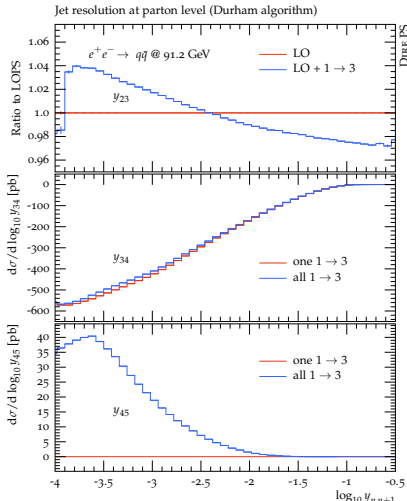
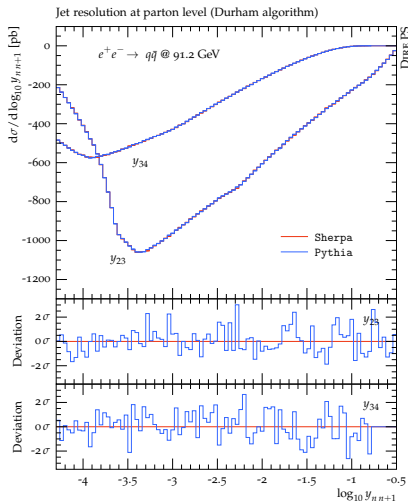
$$C_{qq'}(z) = \int_z \frac{dx}{x} \left( P_{qg}^{(0)}(x) + \varepsilon \mathcal{J}_{qg}^{(1)}(x) \right) \frac{1}{\varepsilon} P_{gq}^{(0)}(z/x)$$

$$\mathcal{J}_{qg}^{(1)}(z) = 2C_F \left( \frac{1 + (1-x)^2}{x} \ln(x(1-x)) + x \right)$$

- ▶ Analytical computation of I not needed, as  $I + \mathcal{P}/\varepsilon$  finite  
generate as endpoint at  $s_{ai} = 0$ , starting from integrand at  $\mathcal{O}(\varepsilon)$
- ▶ All components of  $P_{qq'}^{(1)}$  eventually finite in 4 dimensions  
Can be simulated fully differentially in parton shower

# Higher-order corrections: Collinear evolution at NLO

[Gellersen, Prestel, SH] arXiv:2110.05964



► Effects on jet rates in  $e^+e^- \rightarrow \text{hadrons}$  at LEP

# Higher-order corrections: Multi-Emission Kernels

[Löschner,Plätzer] arXiv:2112.14454

- ▶ Program to define higher-order splitting functions for parton showers
- ▶ Sudakov-like momentum decomposition  $\rightarrow$  power counting
- ▶ Reproduces known soft & double-/triple-collinear splitting functions

$$\begin{aligned}
 & \frac{\mu^{2\epsilon}}{\hat{\alpha}^2 S_{i12}^2} \left\{ \begin{array}{l} \text{Diagram 1} + \text{Diagram 2} + \\ \text{Diagram 3} + \text{Diagram 4} + (1 \leftrightarrow 2) \end{array} \right\} C_A C_F \\
 & = \left( \frac{8\pi\alpha_S}{\hat{\alpha} S_{i12}} \mu^\epsilon \right)^2 C_A C_F \langle \hat{P}_{ggq}^{(\text{non-Ab})} \rangle \hat{p}_i + \mathcal{O}(\beta_{il}^{-3/2}).
 \end{aligned}$$

# Looking beyond logarithmic accuracy

- ▶ Provably NLL accurate parton showers solve long-standing problem NNLL seems on the horizon, but is it the obvious target?
- ▶ Revisit well-established result: Thrust or  $FC_{1-\beta}$  in  $e^+e^- \rightarrow \text{hadrons}$
- ▶ Define a shower evolution variable  $\xi = k_T^2/(1-z)$
- ▶ Parton-shower one-emission probability for  $\xi > Q^2\tau$

$$R_{\text{PS}}(\tau) = 2 \int_{Q^2\tau}^{Q^2} \frac{d\xi}{\xi} \int_{z_{\min}}^{z_{\max}} dz \frac{\alpha_s(k_T^2)}{2\pi} C_F \left[ \frac{2}{1-z} - (1+z) \right] \Theta(\eta)$$

- ▶ Approximate to NLL accuracy

$$R_{\text{NLL}}(\tau) = 2 \int_{Q^2\tau}^{Q^2} \frac{d\xi}{\xi} \left[ \int_0^1 dz \frac{\alpha_s(k_T^2)}{2\pi} \frac{2C_F}{1-z} \Theta(\eta) - \frac{\alpha_s(\xi)}{\pi} C_F B_q \right]$$

## Origin of the $\alpha_s \rightarrow 0 / s \rightarrow \infty$ limit

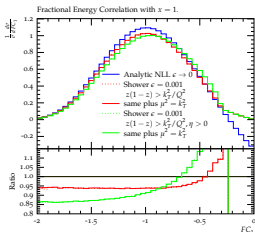
- ▶ Cumulative cross section  $\Sigma(\tau) = e^{-R(\tau)} \mathcal{F}(\tau)$  obtained from all-orders resummed result by Taylor expansion of virtual corrections in cutoff  $\varepsilon$

$$\mathcal{F}(\tau) = \int d^3 k_1 |M(k_1)|^2 e^{-R' \ln \frac{\tau}{\varepsilon v_1}} \sum_{m=0}^{\infty} \frac{1}{m!} \left( \prod_{i=2}^{m+1} \int_{\varepsilon v_1}^{v_1} d^3 k_i |M(k_i)|^2 \right) \\ \times \Theta(\tau - V(\{p\}, k_1, \dots, k_n))$$

- ▶  $\mathcal{F}(\tau)$  is pure NLL & accounts for (correlated) multiple-emission effects
- ▶ In order to make  $\mathcal{F}(\tau)$  calculable, make the following assumptions
  - ▶ Observable is recursively infrared and collinear safe
  - ▶ Hold  $\alpha_s(Q^2) \ln \tau$  fixed, while taking limit  $\tau \rightarrow 0$ 
    - Can factorize integrals and neglect kinematic edge effects
- ▶ Breaks momentum conservation and unitarity for finite  $\tau$ 
  - Clean NLL result, but unknown kinematic corrections
- ▶ How large are effects in regions of a typical measurement?

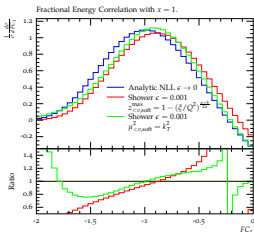
# Numerical effects away from the limit

[Reichelt,Siegert,SH] arXiv:1711.03497



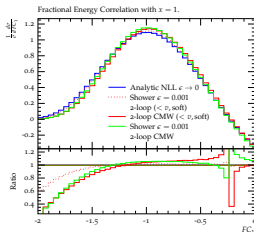
Single emission effects

- ▶ 4-mom conservation
- ▶ PS sectorization
- ▶  $k_T$  scale in coll. terms



Multiple emission effects


- ▶  $z$  bounds by unitarity
- ▶  $k_T$  scale by unitarity



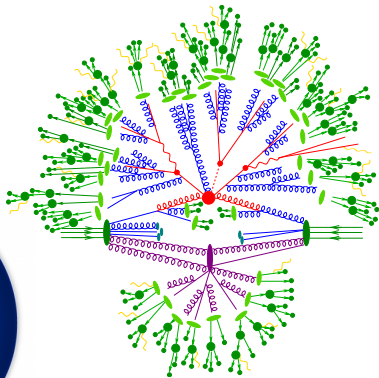
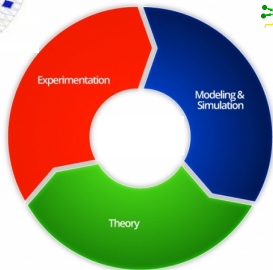
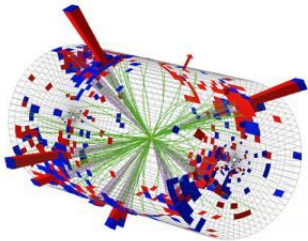
Effects of scale choice

- ▶ Simplest process and simplest type of observable, still sizable differences away from  $\tau \rightarrow 0$  limit
- ▶ How do we quantify the precision of event generators in the intermediate region (“between” NLL and NLO) ?

# Summary and Outlook

- ▶ Lots of activity in event generator development ...
  - ▶ Logarithmic precision of parton showers [PanScales,Herwig,Sherpa,...]
  - ▶ Higher-order QCD evolution kernels [Vincia,Sherpa,Herwig,...]
  - ▶ Interplay of parton showers w/ NNLL [PanScales,Sherpa,...]
  - ▶ Improved & alternative hadronization models [ talk by T. Menzo]
- ▶ ... and matching to fixed-order calculations
  - ▶ Novel computing techniques [MadGraph5,Sherpa]
  - ▶ Resummation based NNLO matching [Geneva,MINNLOps]
  - ▶ Fully differential (N)NNLO matching [Vincia,UN<sup>X</sup>LOPS,TOMTE]
- ▶ Still, many improvements needed [Campbell et al.] arXiv:2203.11110
  - ▶ Systematic treatment of kinematic edge effects
  - ▶ Massive quark production & evolution
  - ▶ Other exciting areas:  $\nu$ s, HI, EIC, ...
  - ▶ ...

**Exciting times ahead!**



$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\Psi}\not{D}\Psi + h.c.$$